## ENABLING REPEAT-PASS INTERFEROMETRY FROM LOW VENUS ORBIT

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Repeat-pass interferometry is a powerful technique for determining changes in topography by flying a radar over the terrain two or more times. These overflights must be very close to each other in space. To design and maintain a low Venus orbit that enables this requires the consideration of drag, non-spherical gravity effects, and solar tides. Once the orbit is designed, the spacecraft must be navigated. To do so requires the use of radar-based terrain-relative navigation in addition to the traditional radiometric datatypes. The mission design and navigation to enable repeat-pass interferometry at Venus are described.

#### INTRODUCTION

A deep understanding of solar system evolution and a better understanding of exoplanet habitability is hampered by the unanswered question: Why are Earth and Venus so different? We know that these twin planets formed with similar bulk composition and size. Yet Venus followed a divergent evolutionary path, losing its surface water and becoming hotter than Mercury. How did this happen? The answer has profound implications for the potential for life in the universe and how terrestrial planets become habitable.

Repeat-pass interferometry (RPI) is a powerful radar technique that can detect changes in the topography of a surface with sub-wavelength precision. These topographic changes can be driven by active magmatic and tectonic processes. Synthetic aperture radars onboard aircraft have detected ground motion on the scales of a few millimeters. To perform RPI, the radar must repeat its path over the topography of interest to within a fraction of the critical baseline length. This critical length defines a "tube" through which the spacecraft must fly.

### REPEAT PASS INTERFEROMETRY

Differential radar interferometry<sup>1</sup> is routinely used to measure millimeter-level surface deformation by acquiring radar observations on temporally separated images spanning a time interval over which the surface was deforming. The vector separating the two platform positions when a

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point is imaged is called the baseline vector and the parallel and perpendicular components of the baseline are the projections onto the radar line-of-sight and cross line-of-sight (in blue) directions as shown in Figure 1.

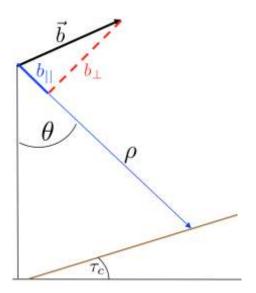


Figure 1: Parallel and perpendicular components of the interferometric baseline

Viable repeat pass interferometry is only possible when the perpendicular component of the baseline is less than the critical baseline,  $b_c$ , given by:

$$b_c = \frac{\lambda \rho \tan(\theta - \tau_c)}{2\Delta \rho} \tag{1}$$

where  $\lambda$  is the radar wavelength,  $\rho$  is the range,  $\theta$  is the look angle,  $\tau_c$  is the cross-track slope and  $\Delta \rho$  is the range resolution of the radar. The critical baseline corresponds to when the geometric correlation,  $\gamma_g$ , given by:

$$\gamma_g = 1 - \frac{b_{\perp}}{b_c} = 1 - \frac{2b_{\perp}\Delta\rho}{\lambda\rho\tan(\theta - \tau_c)}$$
 (2)

becomes zero. We set our minimal allowed geometric correlation to 0.8 to provide for robust repeat pass radar interferometric measurements yielding a maximum tube diameter of 160 meters.

#### MISSION DESIGN

Three perturbations, drag, eccentricity vector evolution due to the non-spherical Venus gravity field, and inclination torqueing due to solar tides drive a Venus orbiter's trajectory away from the Keplerian ideal and an exact repeat. These perturbations must all be considered in the design of the operational orbit and in targeting and navigating to the tube. The design considered here is a representative near-circular, near-polar orbit:  $179 \times 255$  km altitude at 88.5 deg inclination. The perturbations need not necessarily be cancelled in their entirety, but need only be sufficiently cancelled to fly through the tube while the RPI measurements are being taken.

In order to do the interferometry, a full radar dataset must be returned to Earth. The quantity of data and the realities of link budgets across interplanetary distances mean that RPI cannot be done globally. As a result, we consider here "postage-stamp" or "targeted" RPI, where the measurements

are taken for 30 seconds an orbit over 20 orbits, which corresponds, roughly, to a  $200 \times 200$  km patch on the ground. This is the equivalent of looking at Southern California from Santa Barbara to northern San Diego or the San Francisco Bay Area from San Jose to Sacramento and inland to the foothills of the Sierra Nevada.

## **Atmospheric Drag**

The first perturbation is drag, which has been studied at length. The effect of drag is to reduce the semi-major axis of the orbit. For a sufficiently eccentric orbit, it also reduces the eccentricity by lowering apoapsis preferentially. This has two effects on targeting the tube. The altitude change is the first, and obvious, effect. The second is a cross-track effect. If the orbit is lower than the reference, the ground-track will drift to the east. If higher, it will drift west<sup>3</sup>. This leads to the classic "horseshoe" maintenance sequence<sup>4</sup> where the orbit starts above and to the east of the reference altitude and track. Atmospheric drag is then permitted to lower the orbit as it drifts toward the track. Ideally, when the orbit is at the reference altitude, the ground-track has drifted to the edge of its tolerance and the track drift reverses itself as drag continues to lower the orbit. A drag make-up maneuver (or DMU) then raises the orbit when either an altitude or maximum ground-track offset is reached. The initial offsets are typically set such that a lower-than expected drag would still meet the western edge of the ground-track limit, as shown in Figure 2.

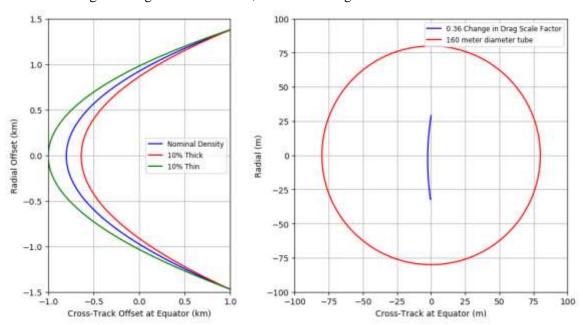


Figure 2: Horseshoe Ground-Track Figure 3: 99<sup>th</sup> Percentile Effect of Drag Uncertainty Maintenance

For the repeat pass interferometry requirement of flying through a 160 meter diameter tube laid down by a previous overflight, the same kind of horseshoe applies. However, the reference is not an absolute altitude and ground-track as in traditional orbit maintenance, but instead what was laid down by the previous trajectory. If the second overflight targets the same altitude and ground-track as was experienced by the first overflight, then the only drag-induced trajectory difference would be due to changes in the atmospheric density over the intervening time. Drag variation is typically parameterized as a scale factor on the atmospheric density model. This covers not just actual density variation but also cross-sectional area, coefficient of drag, and other effects. The 99<sup>th</sup> percentile of variation from one pass to the next is a 0.36 change in scale factor if we conservatively assume that

0.1 is the standard deviation of the scale factor. For an example 218 km circular orbit, a reasonable area/mass ratio, and a scale factor change from 0.829 to 1.197, the second pass will drop 62 meters more than the first pass did in 20 orbits, if the change is compensated for in the targeting, as illustrated in Figure 3. Without compensation, the drop is closer to 72 meters, and the crosstrack difference would grow from 2.3 to 11 meters.

Operationally, the flight team will need to reconstruct and trend the drag scale factors to account for this predominately-radial motion in the tube. If the atmosphere is predicted to be thicker than was experienced in the previous baseline, the trajectory will need to be adjusted to a higher initial condition. The reverse is true if the drag is predicted to be less. The reference altitude will drop faster than the second pass, and so the spacecraft will appear to rise relative to the tube center. Managing the ground-track motion here is far simpler. In both cases, the ground-track will shift a few meters from the reference and need not be specially compensated for, if the altitude is properly compensated. Maintaining the orbit such that the ground-track offset is aimed to be zero at the targeted sites would be an ongoing process.

## **Venus Non-Spherical Gravity**

The second perturbation is due to Venus's extraordinarily slow 243-day sidereal period. The ground-track at the equator only moves about 10 km per orbit. As a result, the same gravity perturbations are experienced repeatedly on subsequent orbits. This drives significant inclination and eccentricity vector evolutions. The inclination varies by  $\pm 0.5$  deg but essentially repeats (see Figures 4 and 5). However, the eccentricity vector sees a secular shift and the altitude vs. ground-track position would be radically different (see Figure 6). A similar effect is experienced by lunar orbiters<sup>5</sup>. The primary effect of this eccentricity vector shift is to force the design away from a circular orbit. To maintain the orbit to within 80 meters of circular would require an eccentricity reset maneuver every few hours. By allowing the orbit to evolve over the course of a ground-track repeat cycle and then resetting at the end of the cycle (or at pre-determined times within a cycle), the operational tempo can be more reasonable. In order to ensure that the spacecraft flies within the tube on subsequent passes, this eccentricity vector evolution must be reset to follow the same path every cycle.

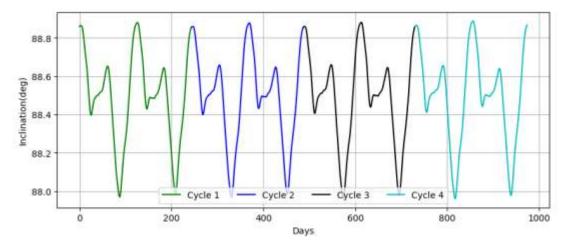


Figure 4: Inclination variation over four ground-track repeat cycles

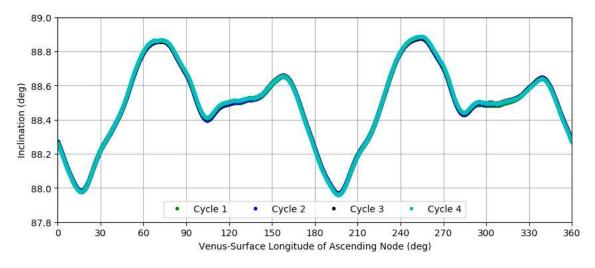


Figure 5: Inclination variation is dominated by Venus non-spherical gravity

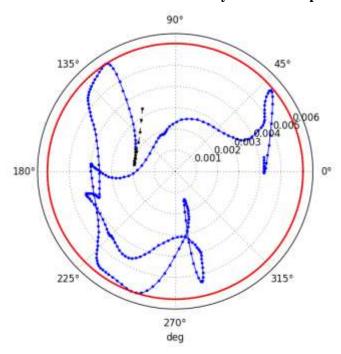


Figure 6: Eccentricity Evolution of a low Venus Orbiter. The first 10 days (black) of one Cycle are not identical to the first 10 days (upper right, in blue) of the following Cycle.

Both of these effects are deterministic (unlike drag) and repeat exactly with the topography (unlike solar torques), and so do not contribute to relative motion of the tube, except insomuch as execution errors in the eccentricity vector resets require clean-up. However, it has been shown<sup>6</sup> that an initial small offset in eccentricity and argument of periapsis (e- $\omega$ ) space (as in Figure 6) of the eccentricity vector remains nearly constant over time. That is, the e- $\omega$  state transition matrix is effectively identity, and residual errors due to clean-ups can be made nearly arbitrarily small if sufficient maneuvers are planned. Small variations in inclination and the ascending node remain similarly small after a 243.3-day ground-track repeat cycle, illustrated in Figure 7. Note that the inclination at the end of a cycle is the same as at the start, but the ascending node has shifted slightly

more than 0.4 degrees. This is the extraordinarily slow precession of the ascending node due to a J2 three order of magnitude smaller than Earth's and is why the repeat cycle is 0.3 days longer than the sidereal day.

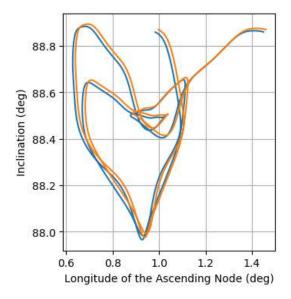


Figure 7: Small variations in inclination and ascending node remain small.

### **Solar Torques**

The third effect is solar torqueing of the orbit plane. If Venus were tidally locked, this would not be an issue, as the solar effect would repeat along with the non-solar gravity effect. However, Venus is not tidally locked; its orbit period is 224 days. The solar perturbation on inclination, averaged over an orbit, is given by:

$$\frac{di}{dt}\Big|_{AV} = -\frac{3\mu_s}{2r_s^3} \sqrt{\frac{a^3}{\mu}} \sin i \sin(\lambda - \Omega) \cos(\lambda - \Omega)$$
 (3)

where  $\mu_s$  and  $\mu$  are the gravitational parameters of the sun and Venus, respectively,  $r_s$  is the range to the sun, a is the orbit semi-major axis, i is the inclination relative to Venus orbital plane,  $\lambda$  is the longitude of the sun, and  $\Omega$  is the orbital ascending node. Venus is in a very nearly circular 0.718  $\times$  0.728 AU orbit with an orbit obliquity of 2.6 deg. From perihelion to aphelion, a low Venus orbiter would only experience a 4% change in this rate. A polar orbit would experience an even smaller rate change due to the obliquity: 0.1%. For conservatism, we shall assume the maximum rate for our 179  $\times$  255 km orbit of 4  $\times$  10<sup>-9</sup> deg/sec (including the sine and cosine terms at their maximum value of 0.5). At maximum latitude, this corresponds to 48 meters of total cross-track motion over 20 orbits. Assuming an 88.5 deg inclination orbit, this is reduced to 42 meters at 60 deg latitude. It is a mere 24 meters at 30 deg latitude. Three ground-track repeat cycles into the mission, the inclination change due to solar torque is almost 180 degrees out of phase from the first cycle, as illustrated Figure 8.

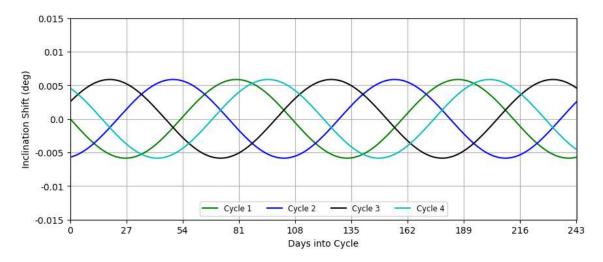


Figure 8: Inclination variation due to solar torque over four ground-track repeat cycles

This  $\pm 0.006$  deg variation in inclination results in a  $\pm 682$  meter crosstrack variation at maximum latitude, far in excess of the  $\pm 80$  meters permissible within the tube, and so this effect must be managed. The operationally simplest way to manage the inclination is to target back to the perturbted inclination of the first observation. That is, make an initial observation of the target in Cycle 1 and then crank the orbit up or down to that inclination in each subsequent cycle. The problem with this approach is that in Cycle 3, the relative inclination rates have almost doubled (see Figure 9), to  $7.8 \times 10^{-9}$  deg/sec, which leads to 94 meters of relative motion. While 94 meters is sufficient to remain in a 160 meter tube with appropriate targeting, it reduces the allowable uncertainty in navigating to the tube entrance. While it costs slightly more propellant to target each observation to an idealized "no-sun" reference as in Figure 8, since doing so requires a maneuver for every target, the reduction in relative tube motion is worth the expense, since a worst-case inclination adjustment is only 0.75 m/s to correct a 0.006 deg offset. The average cost is far less, 0.55 m/s.

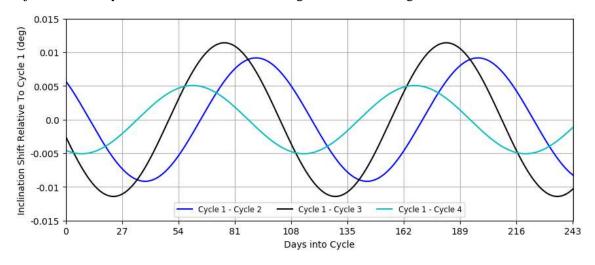


Figure 9: Inclination variation due to solar toques relative to a previous cycle.

## **NAVIGATION**

In general, two-way X-band range-rate tracking data is sufficient for navigation at Venus. However, geometric special conditions arise frequently enough that a combination of two-way Doppler

and two-way-minus-three-way Doppler is a preferable base-line tracking plan. Two-way-minus-three-way Doppler is valuable because it provides a plane-of-sky velocity component – important for reducing cross-track axis error. For brevity, two-way-minus-three-way Doppler will henceforth be referred to as differenced Doppler. Additionally, if the telecommunication subsystem can provide Ka-band downlink, Ka-band tracking data can be base-lined as well. Ka-band adds margin to navigation solutions at Venus, especially for the worst-case geometric condition of an edge-on orbit plane during conjunction.

Daily requirements for navigation are driven by the surface elevation mapping campaign. A successful campaign depends upon on-board processing of the surface measurements – for which an on-board, *predictive* ephemeris is needed. The on-board ephemeris is valid for two to three days and is generated daily from navigation solutions. RPI science requirements, on the other hand, are more stringent in predictive ephemeris with respect to the Venus surface for targeting the tube entry and are driven by ephemeris *reconstruction* accuracy. For the spacecraft to follow a Venus-relative trajectory laid down by earlier orbits over the same site, controllers must be prepared to position two (non-consecutive) orbit paths within 160 m of each other along path at least 200 km in length. Inclination errors, eccentricity errors, orbit-torqueing errors, and maneuver execution errors all contribute to missing the tube. Since Venus's rotation period is uncertain, it too is an error contributor, varying by up to  $\pm 5$  min over decades<sup>7</sup> and perhaps by  $\pm 1$  min over one month<sup>8</sup>.

## **Doppler Performance**

Commanding the spacecraft to thread the tube, once it is known sufficiently well, can be accomplished using the baseline data set of tracking data. Under worst case orbit geometries (orbit plane is within 30 degrees of edge-on toward Earth), and with a 2-day's advance notice prior to tube-entry, the predicted ephemeris cross-track error at tube entry can be as large as 68 m (3 $\sigma$ ) from the target, with components detailed in Table 1. A smaller cross-track contributor is the Venus period error. It yields a two-day prediction error equal to  $\pm 5$  m (for a point on the equator). Radial error for a 2-day prediction is  $\pm 22.5$  m. This includes the traditional and generous margins on the navigation inputs, such as measurement noise and dynamical errors. Under more benign conditions (face-on or nearly face-on orbit geometries) the predicted error at tube-entry varies but is significantly less, approximately  $\pm 50$  m (3 $\sigma$ ).

Table 1: Cross-Track Error Budget for Worst-Case Edge-On Geometry

Cross-Track Error Budget for Repeat-Pass Interferometry	At 30 deg Lat.(m, 3σ)	At 60 deg Lat.(m, 3σ)
Error in knowledge of first-pass orbit in space		
Error in first-pass orbit position with respect to Venus features	±10	±10
Error in Venus feature positions on Venus rotation model	±10	±10
Error in Venus rotation model prediction of Venus rotation	±7.5	±5
Error in control of second-pass orbit position in space		
Error in orbit plane prediction at RPI tube entrance	±65	±65
Error in Venus rotation from timing of spacecraft arrival at tube	±10	±6.5
Error in position from final targeting maneuver execution	±6.5	±10
Total (RSS)	±68	±68

Orbit *reconstruction* of RPI passages can have high fidelity if terrain-relative, radar tie-point data are incorporated to determine the Venus-relative position of the spacecraft (see Table 2). Tie-points are empowering in this application, as discussed in the following section.

#### RADAR-BASED TERRIAN-RELATIVE NAVIGATION

The navigation subsystem can supplement Doppler measurements with radar tie-point observations<sup>9</sup> to compensate for Venus' rotation uncertainty. A radar tie point is the identification of the same surface feature in two or more radar images. Tie points are used to build a frame-tie, relating the J2000 frame to the Venus Body Fixed frame. This precise knitting-together of coordinate frames enables RPI science observations.

#### Data

Radar tie-point observables are range and relative-velocity between the orbiter and the landmark. They are defined by a set of seven data elements, each set specific to a particular orbit and radar swath. See the Appendix for a description.

For RPI reconstruction, terrain-relative observations and two-way Doppler measurements are sufficient to meet science goals. Differenced-Doppler is unnecessary. Thus, two-way X/Ka Doppler combined with terrain-relative measurements satisfy science goals. (Differenced-Doppler is needed for some predictive ephemerides but not for reconstructed ephemerides.) This offers more flexibility for DSN schedulers, as the contingency of planning for station view-period overlaps is eliminated.

We note that the value of terrain-relative observations for orbit determination strengthens as more tie-points are acquired. Fewer than 10 tie-points connecting orbits i and j are ineffective; we suggest a bare minimum of 20 for any pair of orbits. In this preliminary analysis, we simulated 49 tie-points. Reconstructions are even more robust when hundreds of tie-points are identified.

To simulate tie-point observables for this analysis, two spacecraft trajectories were generated: a nominal ephemeris and a 1- $\sigma$  perturbed ephemeris. The difference in platform position between orbits i and j then defines the 1- $\sigma$  offset between subsequent images of each tie-point. The offset represents our expectation of the actual conditions during data-taking – slightly different observed locations of the same tie-point on subsequent orbits. The best estimate of the tie-point's true coordinates is determined by a least-squares fit of multiple views of the same point.

Results from our first attempt at incorporating simulated Venus-centered landmark data into a solution are shown in Figure 10. We used 49 tie-point landmarks acquired over 15 hours (~9 orbits). Landmark locations were determined using a least-squares fit of those observations. The left figure represents range residuals – the distance from the orbiter to the landmark. The right figure represents range-rate residuals – the velocity of the orbiter with respect to the landmark.

The fits are flat, indicating the simulated data, the dynamic models, and the landmark locations are in rough agreement (subject to loose tolerances). The mean in Figure 10-left is 6 m +/- 1.9 km and Figure 10-right's mean is 1.7 m/s +/- 150 m/s. While the scatter around these averages is large, the statistics are tolerable in light of the tie-point simulation procedure. It is a laborious, multi-step, hands-on process subject to a lot of (subtle) rounding error. The 49 tie-point data set is also *de minimus*. However now that a simulation methodology has been established, in future analyses the procedure can be automated, thereby minimizing systematic errors and simplifying the generation of numerous tie-points.

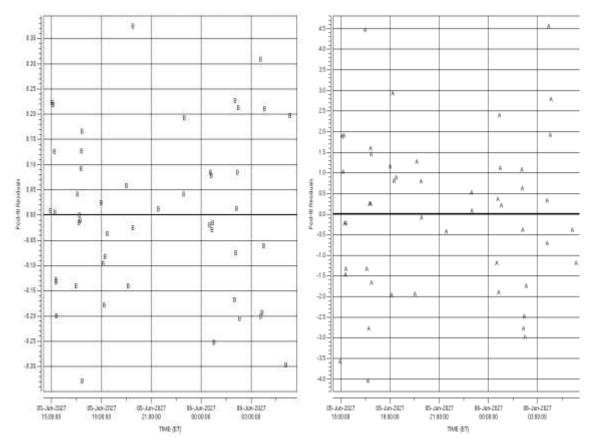


Figure 10: Position (left) and Velocity (right) tie-point residuals for a face-on geometry Reconstruction Results

Combining two-way tracking with terrain-relative measurements reduces cross-track uncertainty significantly with respect to radiometric-only solutions – especially for an edge-on configuration. The cross-track error for face-on orbit geometry shrinks from  $\sim\pm100$  m to  $\sim\pm10$  cm, while the cross-track error for an edge-on orbit geometry is reduced to  $\pm5$  meters. See Figure 10.

Table 2: Orbit Reconstruction Capability (9 hours/day X/Ka Doppler, 2-day data arc)

	Without Tie-points		With Tie-points	
Component	Reqt for Initial Radar Processing	Knowledge (3 $\sigma$ )	Reqt for Fi- nal Radar Processing	Knowledge (3σ)
Radial position (m)	±150	±0.5 <sup>FO</sup>	±2	±0.5 <sup>EO</sup>
Radial velocity (m/s)	±2.1	±0.02 <sup>FO</sup>	-	-
Along-track position (m)	±1500	±250 <sup>FO</sup>	±20	$\pm 5^{\mathrm{FO}}$
Along-track velocity (m/s)	±1.05	±0.001 <sup>FO</sup>	-	-
Cross-track position (m)	±960	±110 <sup>EO*</sup>	±20	$\pm 5^{\mathrm{EO}}$
Cross-track velocity (m/s)	±1.5	±0.12 <sup>EO*</sup>	-	-

FO – worst case is face-on orbital geometry EO – worst case is edge-on orbital geometry

<sup>\* –</sup> includes differenced-Doppler tracking

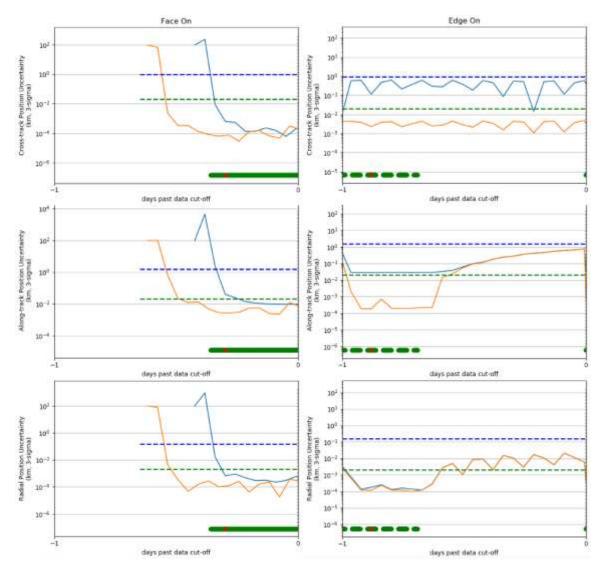


Figure 11: Reconstructed position knowledge for a face-on (left) and edge-on (right) geometries using Doppler tracking only (blue) and with radar tie-points (orange), compared to initial and final radar processing requirements (blue and green dashed lines, respectively)

Table 3 in the Appendix lists modelled error sources for the orbit reconstruction analyses.

## SITE TARGETING

In order to accomplish repeat pass interferometry, each  $200 \times 200$  target must be identified in advance. The ground-track walk would need to be managed continuously such that the offset from the ground-track reference at the target is zero, as in Figure 2. This continuous management would take place within the context of a larger ground-track tolerance set by instrument swath overlap and other, non-RPI considerations. A week prior to the overflight, an inclination adjustment maneuver would be performed to center the projected inclination variation from a no-sun reference, as in Figure 8, on 0 deg variation. Two days prior to the overflight a final maneuver would be performed to achieve the radial target offset by the predicted difference in drag as in Figure 3. Each of these maneuvers can include a component to clean up the execution errors in the previous maneuver(s).

For the first overflight, the radial component is not as critical to achieve, as this will set the reference tube for all future overflights. A worst-case scenario, with an RPI target at 60 deg latitude, with a 99<sup>th</sup> percentile difference in drag, and the highest solar-torque-induced inclination rates combined with the  $2.3\sigma$  (99%)  $104\times35$ -meter delivery ellipse for an edge-on orientation, is shown in Figure 12. The earth is within 30 degrees of edge-on while the inclination rates are within 20% of their maximum value about 28% of the time in one example mission profile, so this geometry is not particularly common, but not exceptionally rare. This is a true worst case and even then the probability of flying within the tube and getting the RPI measurement is very high. Further, if this case were to occur in flight and the delivery errors were to result in a violation at the start (which should happen less than 1% of the time), the relative motion would push it back into the tube and so only the very edge of the  $200\times200$  box would be lost. This edge could be recovered on a third observation when the geometries have changed and either the solar-driven inclination rates would be reduced or the delivery accuracy would improve, or, more likely, both.

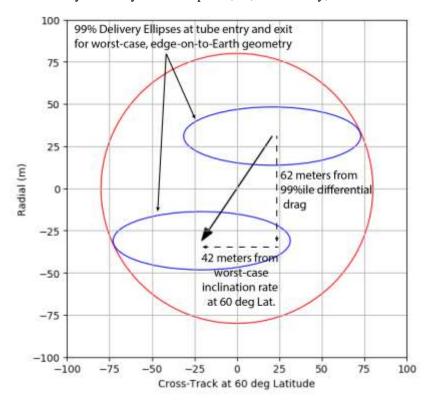


Figure 12: Worst-on-worst relative tube motion at 60 deg latitude (arrow) still meets the 160-m RPI tube (red) with a 99%  $(2.3\sigma)$  delivery (blue) for the worst-case edge-on geometry

### **CONCLUSION**

Repeat pass interferometry is a supremely powerful technique for determining current geologic activity, from the motion of magma chambers deep underground to surface deformations due to earthquakes. It has been used to study these phenomena here on Earth and we can do so at Earth's twin, Venus. By using the enabling technique of radar tie points to knit together the inertial and geocentric coordinate frames of radio-based navigation to the surface of Venus, we can exquisitely reconstruct where the spacecraft previously over-flew the region of interest. This has the by-product of generating a Venus rotation model with an order of magnitude better orientation accuracy than is currently available. Then, using the traditional navigation techniques of Doppler and ranging, we

can target the spacecraft to fly within 80 meters of that previous track with 99% reliability and enable this potentially transformational scientific measurement.

#### **ACKNOWLEDGMENTS**

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### **APPENDIX: RADAR TIE POINTS**

Radar tie points consist of the following information, where i and j denote orbit numbers for two radar swaths, each of which view the same surface feature:

- The times,  $t_j$  and  $t_i$ , when a given feature (or tie point) is imaged by the orbiter on orbits i and j, respectively.
- The Venus Body Fixed positions of the orbiter,  $\sim P_i$  and  $\sim P_j$  when the tie point is imaged on orbits i and j.
- The Venus Body Fixed velocities of the orbiter,  $\sim V_i$  and  $\sim V_j$  when the tie point is imaged on orbits i and j.
- The *a priori* Venus Body Fixed coordinates of the tie point,  $\sim T_i$  and  $\sim T_j$ .
- The ranges  $r_i$  and  $r_j$  when the tie point is imaged on orbits i and j.
- The Doppler frequencies  $f_i$  and  $f_i$  when tie point is imaged on orbits i and j.
- A range-frequency error covariance matrix,  $C_{ij}$ , of the form

$$C_{ij} = \begin{bmatrix} \sigma_r^2 & \sigma_{rf}^2 \\ \sigma_{fr}^2 & \sigma_f^2 \end{bmatrix} \tag{4}$$

**Table 3: Reconstruction Error Sources** 

<b>Error Source</b>	Error $(1\sigma)$	Comments
Orbiter a priori state	30 km (spherical) 100 m/s (spherical)	Wide open initial state.
Landmark location - latitude & longitude	0.5 deg	Notable surface landmarks selected to become radar tie points.
Gravity field (15x15)	MGNP120	GM & covariance from Magellan.
Momentum desaturation impulses	0.6 mm/s (spherical)	Performed once per day.
Tube-targeting $\Delta V$	1.0 cm/s (spherical)	Not part of RPI reconstruction
Tube-targeting cleanup $\Delta V$	5.0 mm/s (spherical)	Not part of RPI reconstruction
Atmosphere density	10%	VenusGRAM2005 atmosphere. Scale factor estimated stochastically. White noise, 8-hour batch.
Solar radiation pressure	5%	Scale factor
UT1, Earth pole $(X,Y)$	21 us, 43 ndeg	Considered

Error Source	Error (1σ)	Comments
Matie (com 0 in)	C4 1 1	Considerati
Media (trop & ion)	Standard error	Considered
DSN station locations	Standard cov.	Considered
Data Type	Measurement error (1σ)	Comments
X/Ka 2-way doppler	4 mhz	= 0.07  mm/s
X/Ka differenced-doppler	3 mhz	= 0.05 mm/s. Scheduled within $\pm 30^{\circ}$ of edge-on orbit geometry
Tie point measurement	1 cm, 50 cm/s	Relative position & velocity of landmark with respect to orbiter at time of radar observation.

#### REFERENCES

<sup>&</sup>lt;sup>1</sup> Rosen, P.A., et al., "Synthetic aperture radar interferometry", Proc. IEEE, Vol. 88, No. 3, March 2000, 333–382

<sup>&</sup>lt;sup>2</sup> Freeman, A., et al, "VERITAS: a Discovery-Class Venus Surface Geology and Geophysics Mission" IEEE Aerospace Conference, Big Sky, MT, March 5-12, 2016

<sup>&</sup>lt;sup>3</sup> Bhat, R.S., et al, "TOPEX/Poseidon Orbit Maintenance for the First Five Years" AAS 98-379 AAS/GSFC 13<sup>th</sup> International Symposium on Space Flight Mechanics, Greenbelt Maryland, May 11-15, 1998

<sup>&</sup>lt;sup>4</sup> Vincent, M. A. "The Mutual Interactions between Drag Make-up and Mean Local Time Maneuver Designs" AIAA/AAS Astrodynamics Specialist Conference, Honolulu, Hawaii, August 18-21, 2008.

<sup>&</sup>lt;sup>5</sup> Sweetser, T. H., et al. "Design of an Extended Mission for GRAIL." AIAA/AAS Astrodynamics Specialist Conference, Minneapolis, Minnesota, August 13-16, 2012.

<sup>&</sup>lt;sup>6</sup> Wallace, M. S. et al., "Low Lunar Orbit Design via Graphical Manipulation of Eccentricity Vector Evolution," AIAA/AAS Astrodynamics Specialist Conference, Minneapolis, Minnesota, August 13-16, 2012.

<sup>&</sup>lt;sup>7</sup> Mueller, N.T., et al. "Rotation period of Venus estimated from Venus Express VIRTIS images and Magellan altimetry". *Icarus* 217, 474–483.doi:10.1016/j.icarus.2011.09.026, 2012

<sup>&</sup>lt;sup>8</sup> Cottereau, L., et al. "The various contributions in Venus rotation rate and LOD." *Astronomy & Astrophysics* 531:A45, 1–10, doi:10.1051/0004-6361/201116606, 2011

<sup>&</sup>lt;sup>9</sup> Chodas, P.W., et al., "Magellan ephemeris improvement using synthetic aperture radar landmark measurements." AAS/AIAA Astrodynamics Conference 1991, Vol. 76 of the *Advances in the Astronautical Sciences* series, 875–890, Univelt, Inc., San Diego, 1992.